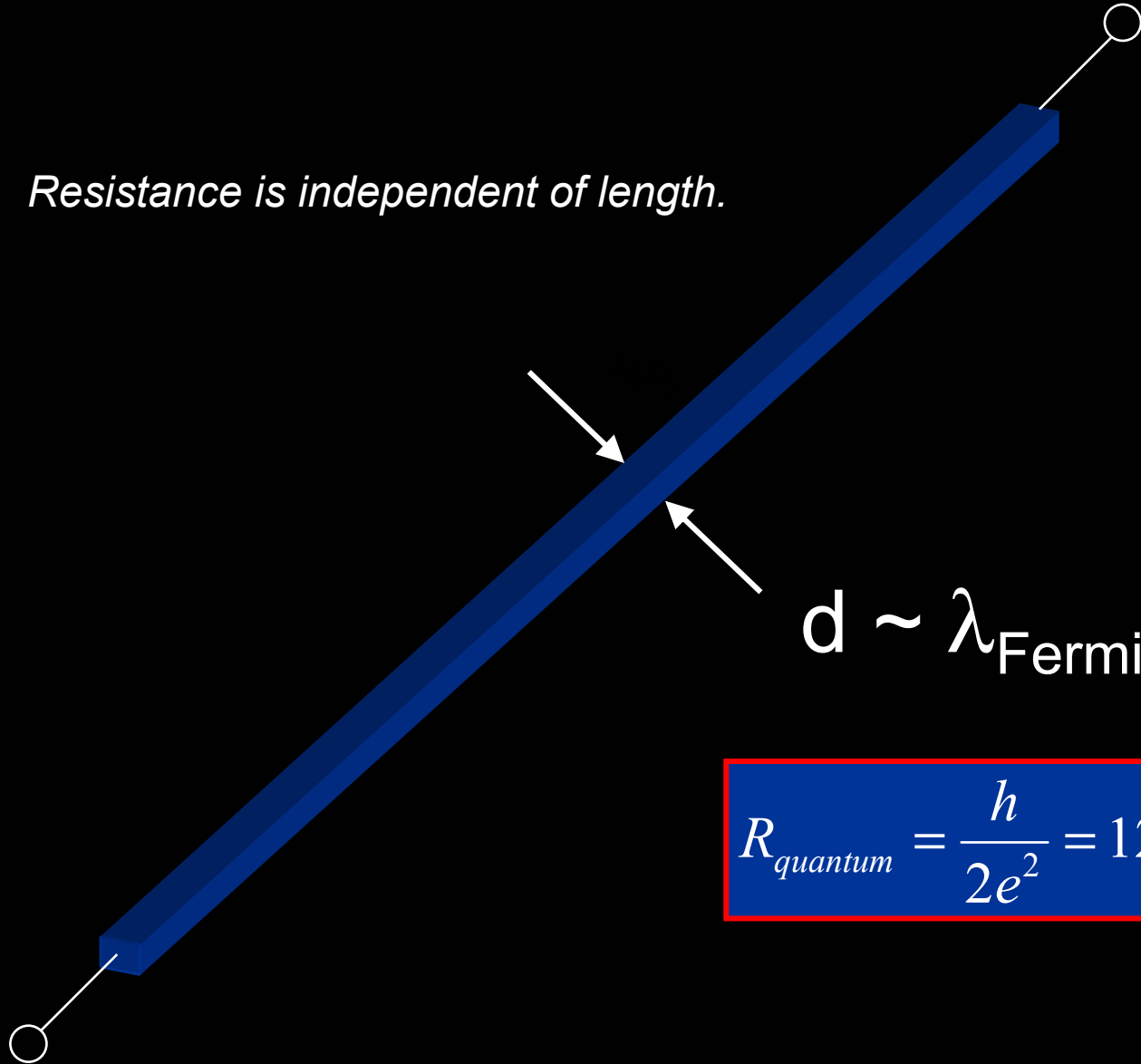


# Lecture 11: Quantum point contact

*Resistance is independent of length.*

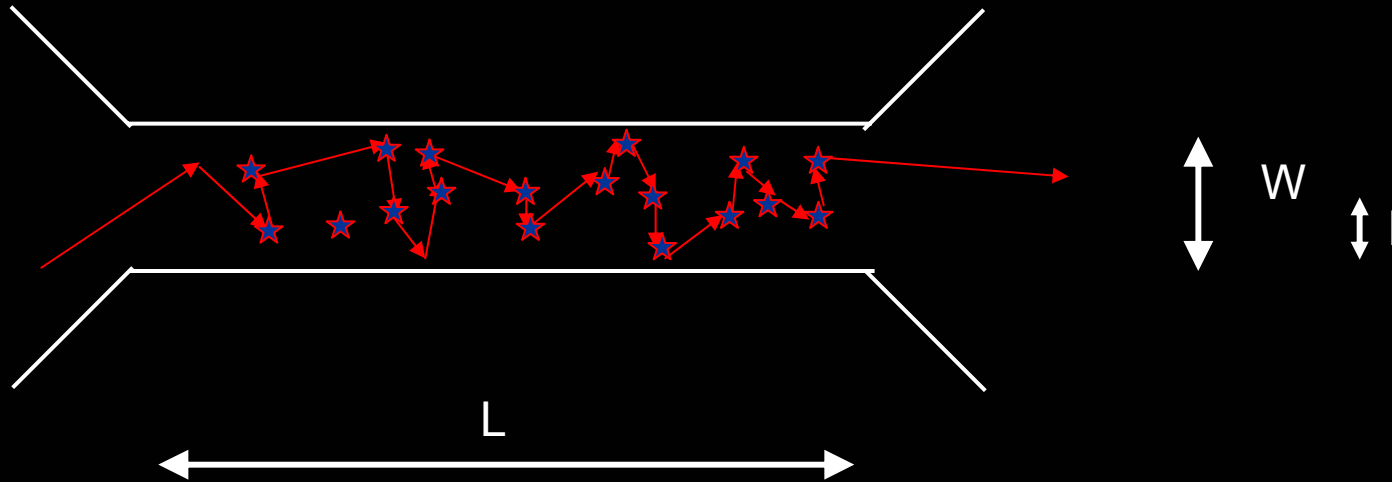


$$R_{\text{quantum}} = \frac{h}{2e^2} = 12.5 \text{ k}\Omega$$

# Ballistic vs. diffusive transport

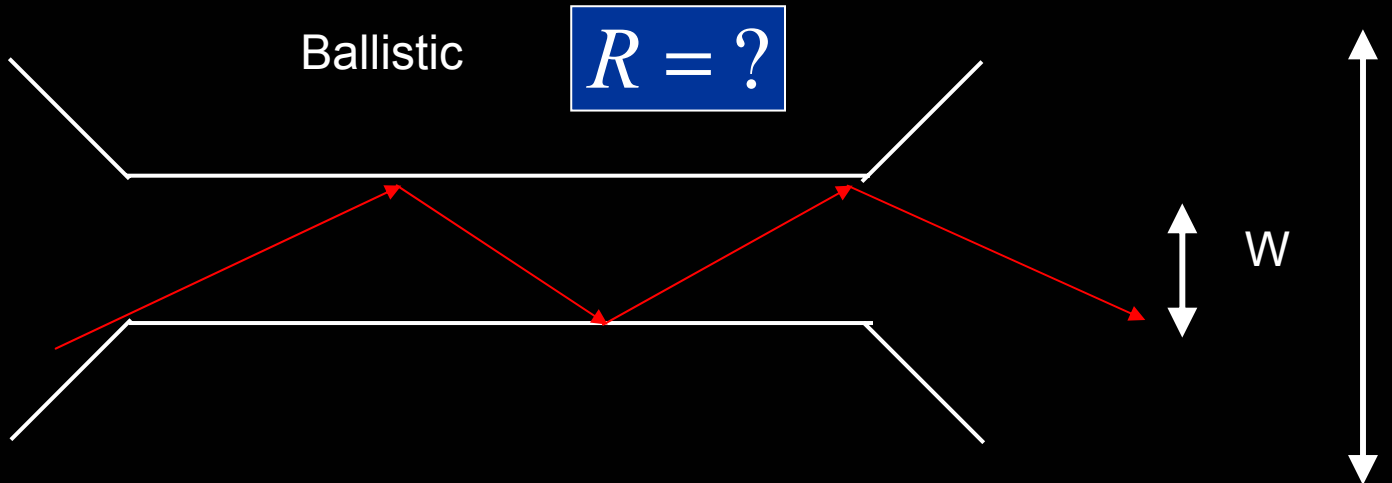
Diffusive

$$R = \frac{L}{W^2} \rho$$



Ballistic

$$R = ?$$



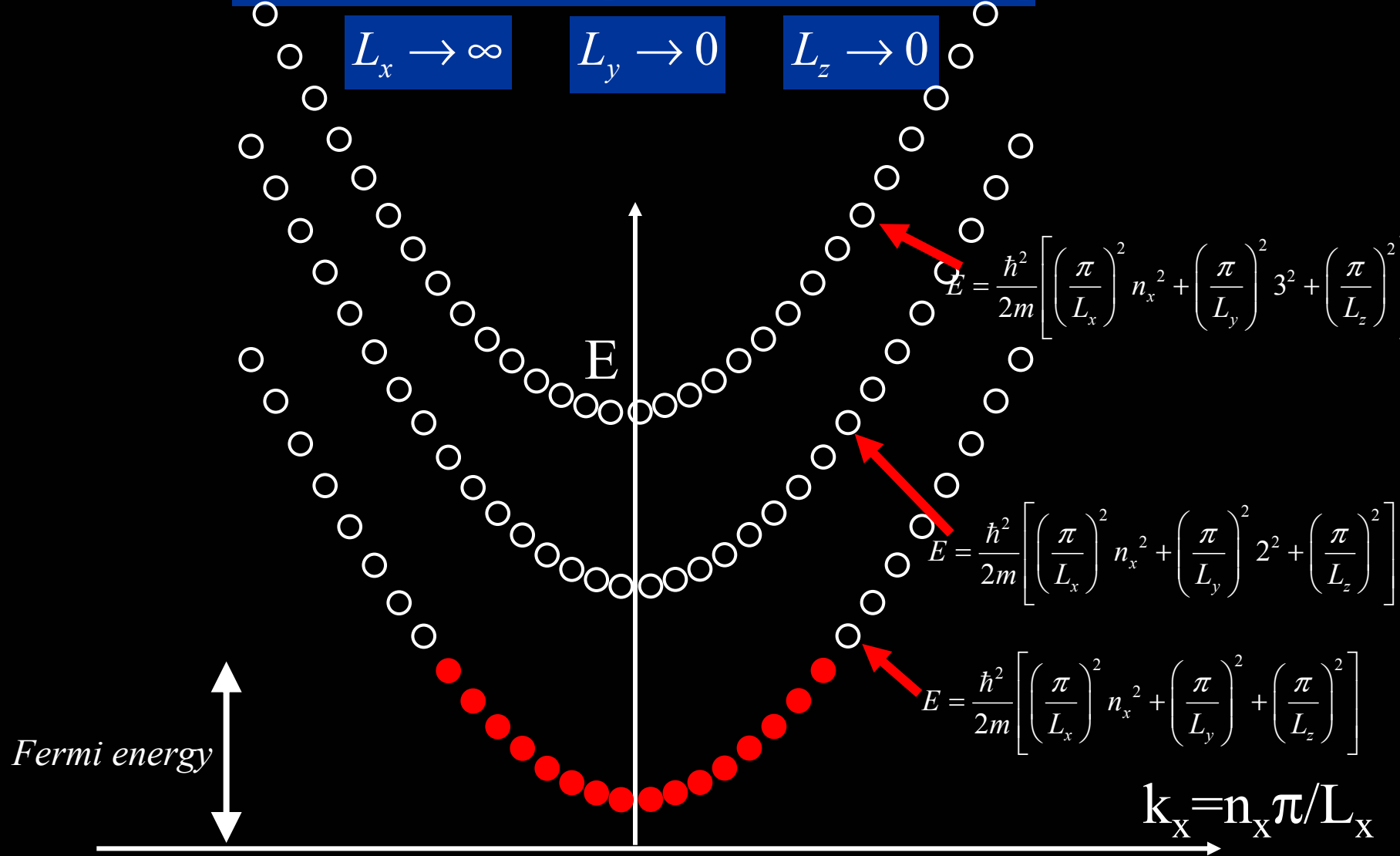
# 1d system:

$$E = \frac{\hbar^2(k_{n_x}^2 + k_{n_y}^2 + k_{n_z}^2)}{2m} = \frac{\hbar^2}{2m} \left[ \left(\frac{\pi}{L_x}\right)^2 n_x^2 + \left(\frac{\pi}{L_y}\right)^2 n_y^2 + \left(\frac{\pi}{L_z}\right)^2 n_z^2 \right]$$

$$L_x \rightarrow \infty$$

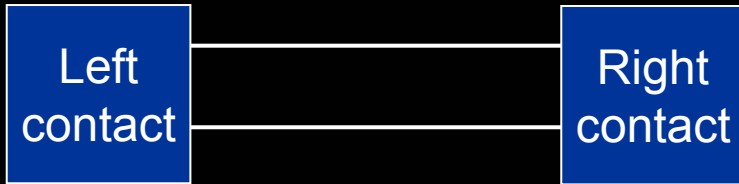
$$L_y \rightarrow 0$$

$$L_z \rightarrow 0$$



# Resistance quantum

Ballistic conductor



$$R_{\text{quantum}} = \frac{h}{e^2} = 25 \text{ k}\Omega$$

With spin:

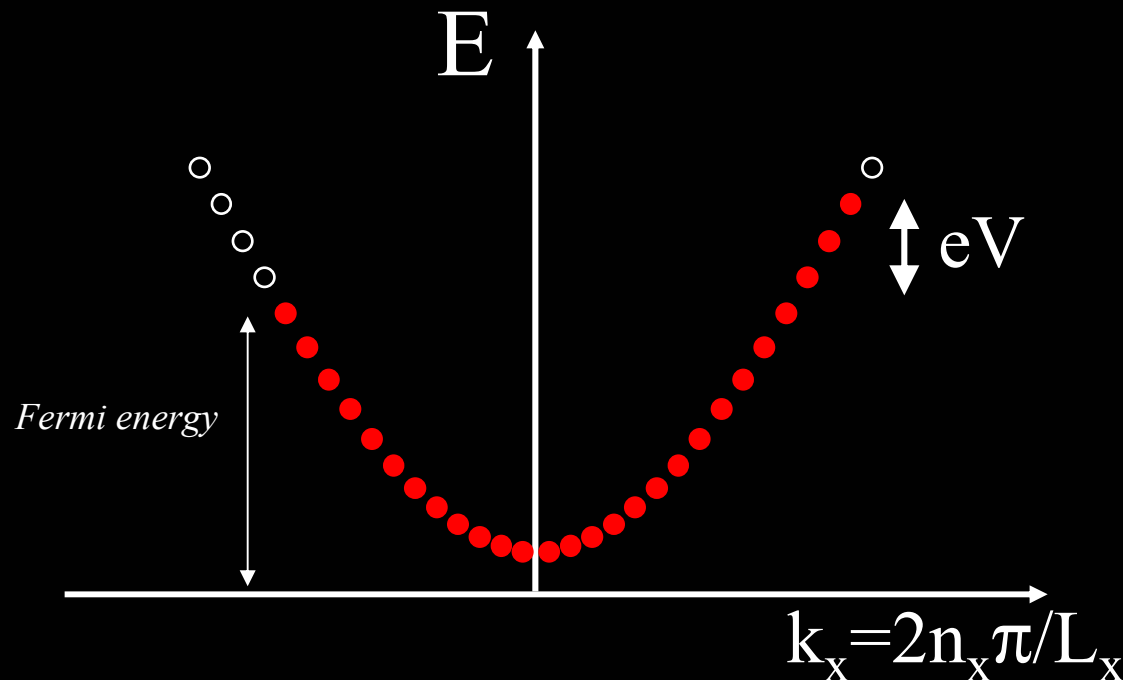
$$R_{\text{quantum}} = \frac{h}{2e^2} = 12.5 \text{ k}\Omega$$

$$G_{\text{quantum}} = \frac{2e^2}{h}$$

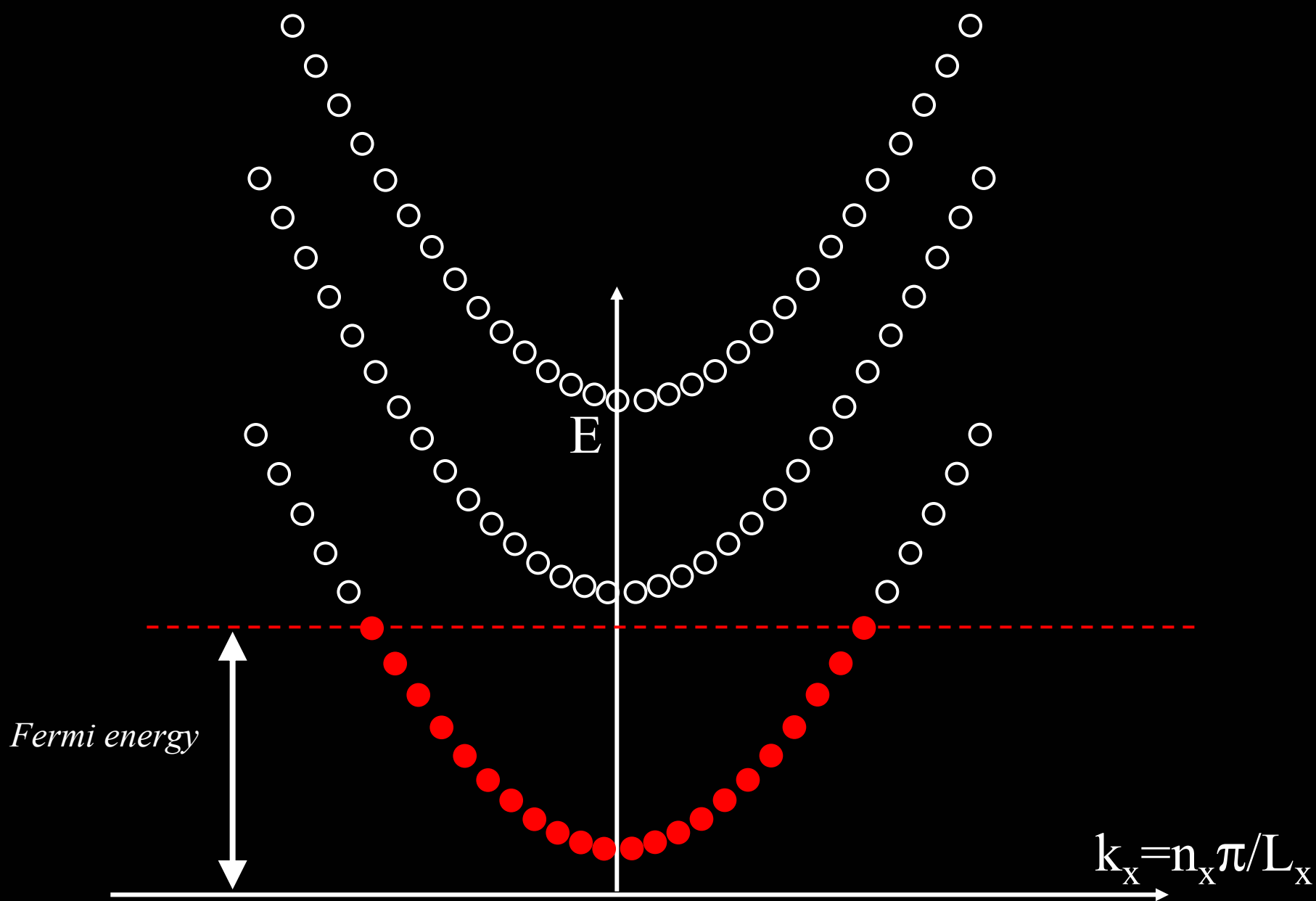
If injection from leads is not perfect:

$$G = T \frac{2e^2}{h}$$

T is the transmission probability.



# Variable width wire:



# Landauer formula:

$$G = n \frac{2e^2}{h}$$

If the leads are not perfect injectors into each “channel” then:

$$G = \frac{2e^2}{h} \sum T_n$$

# Experimental realizations:

- Pinch-off gate in semiconductor 2DEG (QPC)
- Break junction
- Electrochemical addition of atoms
- Scanning tunneling microscope

# Quantum point contact

In class we will show images from:

B.J. van Wees et al. (1988), Phys. Rev. Lett., **60**, 848.

# 0.7 anomaly

In class we will show images from:

<http://marcuslab.harvard.edu/spinfilters.jpg>

# Break junction

In class we will show images from:

Zhou, et al, Applied Physics Letters **67**, 8 (1995) p. 1160.

# Electroplating

In class we will show images from:

A.F.Morpurgo, C.M.Marcus and D. B. Robinson,  
Controlled Fabrication of Metallic Electrodes with Atomic Separation, Appl. Phys. Lett. **74**, 2084 (1999).

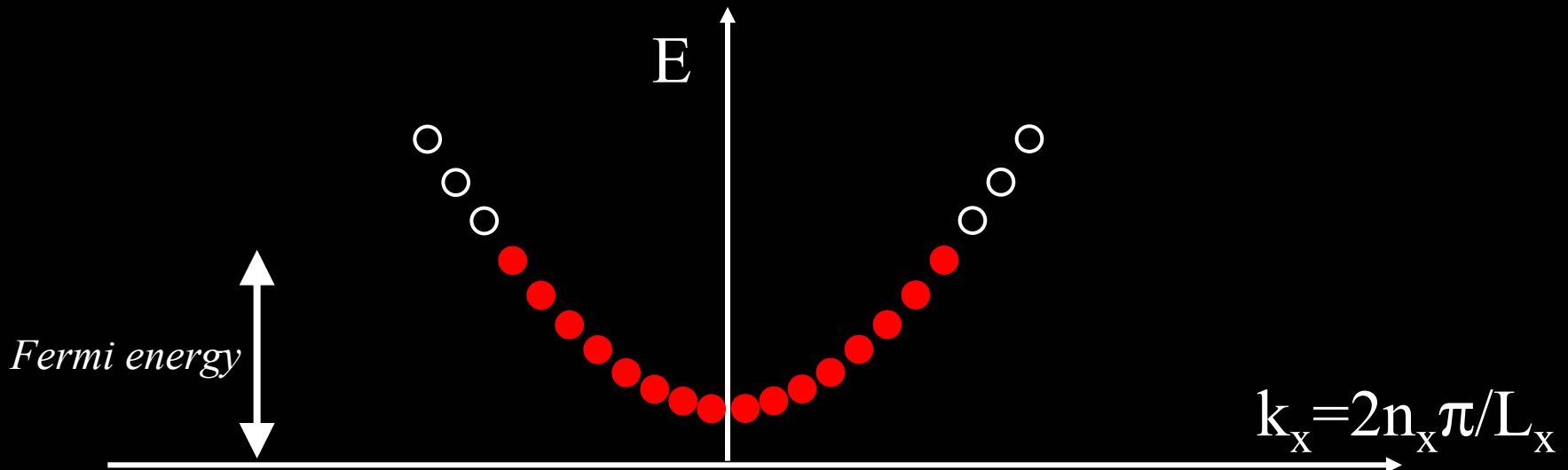
# Positive and negative k-vectors:

Particle in a box: (positive k-vectors only)

$$E = \frac{\hbar^2(k_{n_x}^2 + k_{n_y}^2 + k_{n_z}^2)}{2m} = \frac{\hbar^2}{2m} \left[ \left( \frac{\pi}{L_x} \right)^2 n_x^2 + \left( \frac{\pi}{L_y} \right)^2 n_y^2 + \left( \frac{\pi}{L_z} \right)^2 n_z^2 \right]$$

“Born-Von Karman” boundary conditions: (positive *and* negative k-vectors)

$$E = \frac{\hbar^2}{2m} \left[ \left( \frac{2\pi}{L_x} \right)^2 n_x^2 + \left( \frac{2\pi}{L_y} \right)^2 n_y^2 + \left( \frac{2\pi}{L_z} \right)^2 n_z^2 \right]$$



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- Landauer-Buttiker
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**KeyWords Plus:**  
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*Materials Department, University of California, Santa Barbara, California 93106*

(Received 11 February 1997; accepted 15 April 1997)

The results of high-frequency mixing experiments performed upon parallel quantum point contacts defined in the two-dimensional electron gas of an  $\text{Al}_x\text{Ga}_{1-x}\text{As}/\text{GaAs}$  heterostructure are presented. The parallel geometry, fabricated using a novel double-resist technology, enables the point-contact device to be impedance matched over a wide frequency range and, in addition, increases the power levels of the mixing signal while simultaneously reducing the parasitic source-drain capacitance. Here, we consider two parallel quantum point-contact devices with 155 and 110 point contacts, respectively, both devices operated successfully at liquid helium and liquid nitrogen temperatures with a minimal conversion loss of 13 dB. ©1997 American Institute of Physics.

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